NOTES FROM THE WORKSHOP ON APPROXIMATE GROUPS AND APERIODIC ORDER

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Based on joint work with Michael Björklund

1. Heat kernel hyperuniformity

Let Λ be an invariant (locally square-summable) point process on hyperbolic space

$$\mathbb{H}^n = \mathrm{SO}^+(1, n) / \mathrm{SO}(n)$$

and let σ_{Λ} denote the associated (centered) diffraction measure on $(0, +\infty) \cup i[0, \rho) \subset \mathbb{C}$, where $\rho = (n-1)/2$. Recall from the talk that σ_{Λ} is the unique positive Radon measure on $(0, +\infty) \cup i[0, \rho)$ satisfying

$$\operatorname{Var}\left(\sum_{p\in\Lambda}f(p)\right) = \int_0^\infty |\widehat{f}(\lambda)|^2 d\sigma_{\Lambda}^{(p)}(\lambda) + \int_0^\rho |\widehat{f}(i\lambda)|^2 d\sigma_{\Lambda}^{(c)}(\lambda)$$

for all compactly supported continuous functions f on \mathbb{H}^n , where $\sigma_{\Lambda}^{(p)}, \sigma_{\Lambda}^{(c)}$ are the restrictions of σ_{Λ} to $(0, +\infty)$ and $i[0, \rho)$ respectively.

Assumption: We will assume throughout this note that $\sigma_{\Lambda}^{(c)} = 0$, saving us from any "super-Poissonian" density fluctuations of the process.

Goal: We want show that such a point process Λ is spectrally hyperuniform, i.e.

$$\sigma_{\Lambda}([0,\varepsilon]) = o(\varepsilon^3)$$

if and only if

$$\limsup_{t \to +\infty} t^{3/2} \mathrm{e}^{2\rho^2 t} \operatorname{Var}\left(\sum_{p \in \Lambda} h_t(p)\right) = 0, \qquad (1.1)$$

where

$$h_t(x) = \int_0^\infty e^{-t(\rho^2 + \lambda^2)} \phi_\lambda(d(x, o)) \frac{d\lambda}{|c_n(\lambda)|^2}$$

is the heat kernel on \mathbb{H}^n . If Λ satisfies equation (1.1), then we say that it is *heat kernel hyperuniform*.

1.1. Spectral hyperuniformity implies heat kernel hyperuniformity

Assume that Λ is spectrally hyperuniform. We will first rewrite the variance of the periodized heat kernel into a suitable form and then formulate Lemma 1.1, where the heart of the proof lies.

Consider the function $f_t = e^{\rho^2 t} h_t$ on \mathbb{H}^n , so that

$$\widehat{f}_t(\lambda) = e^{\rho^2 t} \widehat{h}_t(\lambda) = e^{\rho^2 t} e^{-t(\rho^2 + \lambda^2)} = e^{-t\lambda^2} =: \psi(\sqrt{t}\lambda)$$

Using the diffraction measure σ_{Λ} , we rewrite the heat kernel variance as

$$t^{3/2} \mathrm{e}^{2\rho^2 t} \operatorname{Var}\left(\sum_{p \in \Lambda} h_t(p)\right) = t^{3/2} \operatorname{Var}\left(\sum_{p \in \Lambda} f_t(p)\right) = t^{3/2} \int_0^\infty \psi(\sqrt{t}\lambda)^2 d\sigma_\Lambda(\lambda) \,.$$

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Letting $s = \sqrt{t}$, we have to show that

$$s^3 \int_0^\infty \psi(s\lambda)^2 d\sigma_\Lambda(\lambda) \longrightarrow 0 , \quad s \to +\infty .$$

Set

$$\varphi(r) = -2\psi'(r)\psi(r) = 4r e^{-2r^2}$$

so that

$$\psi(s\lambda)^2 = s \int_{\lambda}^{\infty} \varphi(sr) dr$$
.

By a standard change of variables,

$$\int_0^\infty \psi(s\lambda)^2 d\sigma_\Lambda(\lambda) = s \int_0^\infty \left(\int_\lambda^\infty \varphi(sr) dr \right) d\sigma_\Lambda(\lambda) = s \int_0^\infty \varphi(sr) \left(\int_0^t d\sigma_\Lambda(\lambda) \right) dr$$
$$= s \int_0^\infty \varphi(sr) \sigma_\Lambda([0,r]) dr.$$

This means that it is enough to show that

$$s^4 \int_0^\infty \varphi(sr) \sigma_\Lambda([0,r]) dr \longrightarrow 0, \quad s \to +\infty.$$

The core of the proof lies in the following Lemma.

Lemma 1.1 (Björklund). Suppose that there are constants $\alpha, \beta > 0, \gamma \in (0, 1), A, B, C \ge 0$ and Borel functions $\omega, \varphi : [0, +\infty) \to [0, \infty)$ satisfying

(i) ρ is increasing and such that

- for every $\varepsilon > 0$ there is an $r_{\varepsilon} \in (0,1)$ with $r_{\varepsilon} \to 0$ such that $\omega(r) \leq A\varepsilon r^{\alpha}$ for all $r \in [0, r_{\varepsilon}^{1-\gamma}],$
- there is a constant M > 0 such that $\omega(r) \leq Br^{\beta}$ for all $r \geq M$.

(ii) φ is such that

(1) $\int_0^\infty r^\alpha \varphi(r) dr < +\infty,$

(2)
$$R^{\alpha/\gamma} \int_{R}^{\infty} \varphi(r) dr \to 0 \text{ as } R \to +\infty,$$

(3)
$$R^{\alpha-\beta} \int_{R}^{\infty} r^{\beta} \varphi(r) dr \to 0 \text{ as } R \to +\infty.$$

Then

$$r_{\varepsilon}^{-(\alpha+1)} \int_{0}^{\infty} \varphi(r/r_{\varepsilon}) \omega(r) dr \longrightarrow 0 \,, \quad \varepsilon \to 0^{+} \,.$$

With this Lemma in mind, we set $\varphi(r) = 4re^{-2r^2}$ and $\omega(r) = \sigma_{\Lambda}([0, r])$. By the assumed spectral hyperuniformity, there is a constant $A \ge 0$ and a sequence $r_{\varepsilon} \to 0$ such that

$$\omega(r) \le A\varepsilon r^3, \quad r \in [0, r_{\varepsilon}^{1/2}],$$

so we can take $\alpha = 3$, $\gamma = 1/2$. Finding constants B, β as in the Lemma above requires some more work.

Lemma 1.2. Let Λ be an invariant point process on \mathbb{H}^n with diffraction σ_{Λ} . Then $\sigma_{\Lambda}([0, R]) \ll R^n$ for sufficiently large R > 0.

Remark 1.3. When $\omega(r) = \sigma([0, r])$ is the diffraction measure of a point process, Lemma 1.1 is applicable for all rank one spaces. In particular, the crucial constants for Euclidean and hyperbolic space are

G/K	\mathbb{R}^{n}	\mathbb{H}^n
α	n	3
β	n	n

Thus for hyperbolic space there is a B > 0 and an M > 0 such that $\rho(r) \leq Br^n$ for all $r \geq M$, so take $\beta = n$.

Modulo the proof of Lemma 1.1, it suffices to verify points (1)-(3) for the function $\varphi(r) = 4re^{-2r^2}$.

(1) we have

$$4\int_0^\infty r^4 \mathrm{e}^{-2r^2} dr < +\infty \,.$$

(2)

$$R^6 \int_R^\infty 4r \mathrm{e}^{-2r^2} dr = 2R^6 \mathrm{e}^{-2R^2} \longrightarrow 0, \quad R \to +\infty.$$

(3)

$$R^{3-n} \int_{R}^{\infty} r^{n} \mathrm{e}^{-2r^{2}} dr = R^{4} \int_{1}^{\infty} r^{n} \mathrm{e}^{-2R^{2}r^{2}} dr \longrightarrow 0, \quad R \to +\infty$$

by dominated convergence.

Letting $s_{\varepsilon} = r_{\varepsilon}^{-1}$, then Lemma 1.1 tells us that $4\int_{-\infty}^{\infty} f(x) = f(x) dx$

$$s_{\varepsilon}^{4} \int_{0}^{\infty} \varphi(s_{\varepsilon}r) \sigma_{\Lambda}([0,r]) dr \longrightarrow 0 \,, \quad \varepsilon \to 0^{+} \,,$$

which is what we wanted to show.

Proof of Lemma 1.1. Let $I_{\varepsilon} = \int_0^{\infty} \varphi(r/r_{\varepsilon})\rho(r)dr$. Then we use the assumed monotone growth and estimates on ω to find that

$$I_{\varepsilon} \leq \int_{0}^{r_{\varepsilon}^{1-\gamma}} \varphi(r/r_{\varepsilon}) A\varepsilon r^{\alpha} dr + \omega(M) \int_{r_{\varepsilon}^{1-\gamma}}^{M} \varphi(r/r_{\varepsilon}) dr + \int_{M}^{\infty} \varphi(r/r_{\varepsilon}) Br^{\beta} dr =: I_{1} + I_{2} + I_{3}.$$
We need to show that $I_{\varepsilon} = \varepsilon(r^{\alpha+1})$ for $i = 1, 2, 2$ as $\varepsilon \to 0^{+}$

We need to show that $I_j = o(r_{\varepsilon}^{\alpha+1})$ for j = 1, 2, 3 as $\varepsilon \to 0^+$.

$$I_{1} \leq A\varepsilon r_{\varepsilon}^{\alpha+1} \int_{0}^{\infty} \varphi(r) r^{\alpha} dr$$
$$I_{2} \leq \varphi(M) r_{\varepsilon} \int_{r_{\varepsilon}^{-\gamma}}^{\infty} \varphi(r) dr$$
$$I_{3} = B r_{\varepsilon}^{\beta+1} \int_{M/r_{\varepsilon}}^{\infty} \varphi(r) r^{\beta} dr$$

Dividing by $r_{\varepsilon}^{\alpha+1}$, we get that

$$\frac{I_{\varepsilon}}{r_{\varepsilon}^{\alpha+1}} = \frac{I_1 + I_2 + I_3}{r_{\varepsilon}^{\alpha+1}} \ll \varepsilon + \omega(M) \underbrace{(r_{\varepsilon}^{-\gamma})^{\alpha/\gamma} \int_{r_{\varepsilon}^{-\gamma}}^{\infty} \varphi(r) dr}_{\to 0} + \underbrace{r_{\varepsilon}^{-(\alpha-\beta)} \int_{M/r_{\varepsilon}}^{\infty} \varphi(r) r^{\beta} dr}_{\to 0} .$$

1.2. Heat kernel hyperuniformity implies spectral hyperuniformity

Assume that equation (1.1) holds. We will show that $\sigma_{\Lambda}([0, \varepsilon]) = o(\varepsilon^3)$ as $\varepsilon \to 0^+$.

This is fortunately quite straightforward: Let $\delta_t > 0$ and bound canonically from above, $c^{\delta_t} = c^{\delta_t}$

$$\operatorname{Var}\left(\sum_{p\in\Lambda}h_t(p)\right) = \int_0^\infty |\widehat{h}_t(\lambda)|^2 d\sigma_\Lambda(\lambda) \ge \int_0^{\sigma_t} |\widehat{h}_t(\lambda)|^2 d\sigma_\Lambda(\lambda) = e^{-2\rho^2 t} \int_0^{\sigma_t} e^{-2t\lambda^2} d\sigma_\Lambda(\lambda) \,.$$

Assuming heat kernel hyperuniformity, there is for every $\varepsilon > 0$ a $t_{\varepsilon} > 0$ such that

$$t^{3/2} \int_0^{\delta_t} \mathrm{e}^{-2t\lambda^2} d\sigma_{\Lambda}(\lambda) < \varepsilon \,, \quad \forall \, t \ge t_{\varepsilon} \,.$$

In particular,

$$t^{3/2} \mathrm{e}^{-2t\delta_t^2} \sigma_{\Lambda}([0, \delta_t]) < \varepsilon, \quad \forall t \ge t_{\varepsilon},$$

so the question boils down to the following: If we can find $\delta_t \to 0$ as $t \to +\infty$ such that

$$\delta_t^{-3} \ll t^{3/2} \mathrm{e}^{-2t\delta_t^2}, \quad \forall t \ge t_\varepsilon,$$
(1.2)

then

$$\frac{\sigma_{\Lambda}([0,\delta_t])}{\delta_t^3} \ll t^{3/2} \mathrm{e}^{-2t\delta_t^2} \sigma_{\Lambda}([0,\delta_t]) < \varepsilon \,, \quad \forall t \ge t_{\varepsilon} \,,$$

as desired.

Finding a solution δ_t to equation (1.2) is equivalent to $2t\delta_t^2 - 3\log(\delta_t) \leq \frac{3c}{2}\log(t)$ for some c > 0 for all sufficiently large t. As an example, by setting c = 2 we get $2t\delta_t^2 \leq 3\log(t\delta_t)$, for which $\delta_t = t^{-1/2}$ is an example of a solution whenever $t \geq e^{4/3}$.